

A Note on the Degasperis-Procesi Equation

Octavian G MUSTAFA

*Permanent address: Department of Mathematics,
University of Craiova, Al. I. Cuza 13, Craiova, Romania
E-mail: octavian@yahoo.com*

*Present address: Department of Mathematics, Lund University,
P.O. Box 118, SE-22100 Lund, Sweden
E-mail: octavian@yahoo.com*

Received July 09, 2004; Accepted September 13, 2004

Abstract

We prove that smooth solutions of the Degasperis-Procesi equation have infinite propagation speed: they lose instantly the property of having compact support.

1 Introduction

The Degasperis-Procesi equation [13]

$$u_t - u_{txx} + 4uu_x = 3u_x u_{xx} + uu_{xxx}, \quad x \in \mathbb{R}, t \geq 0, \quad (1.1)$$

was derived recently as a shallow water approximation to the Euler equation [14]. The equation (1.1) presents some similarities to the Camassa-Holm equation [1, 17]

$$u_t - u_{txx} + 3uu_x = 2u_x u_{xx} + uu_{xxx}, \quad x \in \mathbb{R}, t \geq 0. \quad (1.2)$$

Indeed [1, 13, 15], both equations are bi-Hamiltonian and have an associated isospectral problem. Therefore they are both formally integrable (the integrability of (1.2) by means of the scattering/inverse scattering approach is discussed in [5, 9, 19]). Also, both equations admit exact peaked solitary wave solutions which have to be understood as weak solutions [10, 8, 22]. Moreover, using Kato's semigroup theory for quasilinear equations of evolution [18], the local existence of solutions to (1.1) in $H^s(\mathbb{R})$, with $s > \frac{3}{2}$ can be established provided the initial data is in $H^s(\mathbb{R})$ (see [23]), and finite-time blowup is possible only if the slope of the solution becomes unbounded in finite time [22]. The last feature parallels the fact that for (1.2) an initial data with the same regularity will either develop into a wave that exists for all times or wave breaking occurs [2, 3, 4, 12].

Despite these similarities, the equations (1.1) and (1.2) are truly different. For example, for (1.1) the isospectral problem is of third order, whereas in the case of (1.2) we encounter a second order isospectral problem [22]. Moreover, (1.2) is a re-expression of geodesic flow on the group of diffeomorphisms of the line [20] (see also [4, 6, 7]) whereas (1.1) does not have a geometric derivation of this type.

A different inquiry is the aim of the present letter. We will prove that the infinite propagation speed of smooth solutions, observed in [11] for the Camassa-Holm equation, is also a feature of the model (1.1).

2 Main results

Denoting $m = u - u_{xx}$, we can write (1.1) in the form

$$m_t + um_x + 3u_x m = 0, \quad x \in \mathbb{R}, t \geq 0. \quad (2.1)$$

Theorem 1. *Let $m_0 : \mathbb{R} \rightarrow \mathbb{R}$ be a smooth function with compact support. If the solution $m(x, t)$ of equation (1.1) has $[0, T)$ as maximal interval of existence in the future, then, at every moment $t \in [0, T)$, the smooth function $m(\cdot, t)$ has compact support.*

Proof. We shall employ a device from [4]. Let us introduce the initial value problem

$$\begin{cases} \frac{d\psi}{dt} = u(\psi, t), & t \in [0, T), \\ \psi(0) = x. \end{cases} \quad (2.2)$$

Since $u_0 = p * m_0$, where $u(x, 0) = u_0(x)$ for $x \in \mathbb{R}$, the local regularity theory developed in [23, 22] for (1.1) implies that $u(t, \cdot)$ is smooth on $[0, T)$. Therefore, via standard qualitative theory for ordinary differential equations [16], the existence in $[0, T)$, uniqueness and smooth dependence of the data x for the solution $\psi = \psi(t; x)$ of (2.2) is ensured. By integration in $[0, t)$, with $t < T$, we get

$$\psi(t; x) = x + \int_0^t u(\psi(s; x), s) ds,$$

which yields

$$\psi_x(t; x) = 1 + \int_0^t u_x(\psi(s; x), s) \psi_x(s; x) ds, \quad t \in [0, T), \quad (2.3)$$

and respectively

$$\frac{d\eta}{dt} = u_x(\psi(t; x), t) \eta,$$

where we have denoted by η the left-hand member of (2.3). Finally, since $\eta(0) = 1$, we obtain

$$\psi_x(t; x) = \eta(t) = \exp\left(\int_0^t u_x(\psi(s; x), s) ds\right), \quad t \in [0, T). \quad (2.4)$$

By introducing the function $\varphi(x, t) = \psi(t; x)$, we have obtained an element of $C^1(\mathbb{R} \times [0, T), \mathbb{R})$. Furthermore, since, due to the Sobolev imbeddings, $u_x(\cdot, t)$ is bounded in \mathbb{R} for all $t \in [0, T)$, we deduce that

$$0 < C_1(t) \leq \psi_x(t; x) \leq C_2(t) < +\infty, \quad x \in \mathbb{R}, t \in [0, T), \quad (2.5)$$

which allows us to conclude that $\varphi(\cdot, t)$, $t \in [0, T)$, are all diffeomorphisms of \mathbb{R} . Also, from (2.2) we get $\varphi_{xt} = u_x \varphi_x$. Further, let us multiply the equation (2.1) by φ_x^3 . Then, if we take $\varphi(x, t)$ and t as arguments of u, m instead of the usual x, t , the result of multiplication reads as

$$\begin{aligned} 0 &= (m_t + um_x + 3u_x m) \varphi_x^3 = m_t \varphi_x^3 + 3m (u_x \varphi_x) \varphi_x^2 + um_x \varphi_x^3 \\ &= m_t \varphi_x^3 + 3m \varphi_{xt} \varphi_x^2 + \varphi_t m_x \varphi_x^3 = \frac{d}{dt} (m \varphi_x^3). \end{aligned}$$

Via an integration,

$$m(\varphi(x, t), t) \varphi_x^3(x, t) = m(\varphi(x, 0), 0) \varphi_x^3(x, 0) = m(x, 0) = m_0(x).$$

Finally, due to (2.5), if the support of m_0 is included in $[a, b]$ then the support of $m(\cdot, t)$ will be included in $[\varphi(a, t), \varphi(b, t)]$. The proof is complete. \blacksquare

Theorem 2. *Let $u_0 : \mathbb{R} \rightarrow \mathbb{R}$ is a smooth function with compact support. If the solution $u(x, t)$ with initial data $u_0(x)$ of (1.1) exists on some time interval $[0, \epsilon)$ with $\epsilon > 0$ and, at every instant $t \in [0, \epsilon)$, the function $u(\cdot, t)$ has compact support, then u is identically zero.*

Proof. According to the preceding theorem, $m(\cdot, t)$ has compact support at every moment $t \in [0, \epsilon)$. We recall that, by the Paley-Wiener theorem [21], an entire (analytic) function $g(\xi)$, where $\xi = \eta + i\zeta$ and $\eta, \zeta \in \mathbb{R}$, is the Fourier transform of a smooth function $f : \mathbb{R} \rightarrow \mathbb{R}$ with compact support in $[-a, a]$ (for $a > 0$), namely

$$g(\xi) = \mathcal{F}_f(\xi) = \int_{\mathbb{R}} f(q) e^{-i\xi q} dq,$$

if and only if for every integer $n \geq 0$ there exists $c_n > 0$ such that

$$|g(\xi)| \leq \frac{c_n e^{a|\zeta|}}{(1 + |\xi|^n)}, \quad \xi \in \mathbb{C}.$$

Since

$$\mathcal{F}_{m(\cdot, t)}(\xi) = (1 + \xi^2) \mathcal{F}_{u(\cdot, t)}(\xi), \quad \xi \in \mathbb{C}, \quad t \in [0, \epsilon),$$

it is obvious that the analiticity of $\mathcal{F}_{u(\cdot, t)}$, if assumed, will imply the analiticity of $\mathcal{F}_{m(\cdot, t)}(\xi)$. In such a case, the function $\mathcal{F}_{m(\cdot, t)}$ has value zero at $i, -i$ for all $t \in [0, \epsilon)$, yielding

$$\int_{\mathbb{R}} e^x m(x, t) dx = \int_{\mathbb{R}} e^{-x} m(x, t) dx = 0, \quad t \in [0, \epsilon).$$

Since both u and m have compact support, we deduce that

$$\begin{aligned}
\frac{d}{dt} \int_{\mathbb{R}} e^x m(x, t) dx &= \int_{\mathbb{R}} e^x m_t dx = -3 \int_{\mathbb{R}} e^x u_x m dx - \int_{\mathbb{R}} e^x m_x u dx \\
&= -3 \int_{\mathbb{R}} e^x u_x m dx + \int_{\mathbb{R}} e^x u_x m dx + \int_{\mathbb{R}} e^x m u dx = -2 \int_{\mathbb{R}} e^x u_x m dx + \int_{\mathbb{R}} e^x m u dx \\
&= -2 \int_{\mathbb{R}} e^x u_x u dx + 2 \int_{\mathbb{R}} e^x u_x u_{xx} dx + \int_{\mathbb{R}} e^x u^2 dx - \int_{\mathbb{R}} e^x u_{xx} u dx \\
&= \left(- \int_{\mathbb{R}} e^x u_x u dx + \frac{1}{2} \int_{\mathbb{R}} e^x u^2 dx \right) - \int_{\mathbb{R}} e^x u_x^2 dx + \int_{\mathbb{R}} e^x u^2 dx \\
&\quad + \left(\int_{\mathbb{R}} e^x u_x^2 dx + \int_{\mathbb{R}} e^x u_x u dx \right) = \frac{3}{2} \int_{\mathbb{R}} e^x u^2 dx = 0,
\end{aligned}$$

which implies that $u(\cdot, t) \equiv 0$ for all t . The proof is complete. \blacksquare

Acknowledgement. This note was written while the author was a Visiting Researcher at Lund University financed by the National Science Foundation of Sweden, Grant VR 621-2003-5287. He is deeply indebted to Professor Adrian Constantin for his warm friendship and continuous support. The author thanks the Mathematics Department at Lund University for making his staying very agreeable.

References

- [1] Camassa R and Holm D, An integrable shallow water equation with peaked solitons, *Phys. Rev. Lett.* **71** (1993), 1661-1664.
- [2] Constantin A and Escher J, Wave breaking for nonlinear nonlocal shallow water equations, *Acta Math.* **181** (1998), 229-243.
- [3] Constantin A and Escher J, Well-posedness, global existence, and blowup phenomena for a periodic quasi-linear hyperbolic equation, *Comm. Pure Appl. Math.* **51** (1998), 475-504.
- [4] Constantin A, Existence of permanent and breaking waves for a shallow water equation: a geometric approach, *Ann. Inst. Fourier (Grenoble)* **50** (2000), 321-362.
- [5] Constantin A, On the scattering problem for the Camassa-Holm equation, *Proc. R. Soc. London Ser. A - Math. Phys. Eng. Sci.*, **457** (2001), 953-970.
- [6] Constantin A and Kolev B, On the geometric approach to the motion of inertial mechanical systems, *J. Phys. A* **35** (2002), R51-R79.
- [7] Constantin A and Kolev B, Geodesic flow on the diffeomorphism group of the circle, *Comment. Math. Helv.* **78** (2003), 787-804.
- [8] Constantin A and Strauss W, Stability of peakons, *Comm. Pure Appl. Math.* **53** (2000), 603-610.
- [9] Constantin A and McKean H P, A shallow water equation on the circle, *Comm. Pure Appl. Math.* **52** (1999), 949-982.
- [10] Constantin A and Molinet L, Global weak solutions for a shallow water equation, *Comm. Math. Phys.* **211** (2000), 45-61.

-
- [11] Constantin A, Finite propagation speed for the Camassa-Holm equation, *Preprint*, Lund University, Sweden, 2004.
- [12] Danchin R, A few remarks on the Camassa-Holm equation, *Diff. Integr. Eqs.* **14** (2001), 953–988.
- [13] Degasperis A and Procesi M, Asymptotic integrability, in: A. Degasperis, G. Gaeta (Eds.), *Symmetry and Perturbation Theory (SPT 98)*, Rome, December 1998, World Scientific, River Edge, NJ, 1999, 23-37.
- [14] Dullin H R, Gottwald G A and Holm D D, Camassa-Holm, Korteweg-de Vries-5 and other asymptotically equivalent equations for shallow water waves, *Fluid Dynam. Res.* **33** (2003), 73-95.
- [15] Fuchssteiner B and Fokas A S, Symplectic structures, their Bäcklund transformation and hereditary symmetries, *Phys. D* **4** (1981), 47–66.
- [16] Hartman P, Ordinary Differential Equations, J. Wiley & Sons, New York-London-Sydney, 1964.
- [17] Johnson R S, Camassa-Holm, Korteweg-de Vries and related models for water waves, *J. Fluid Mech.* **455** (2002), 63-82.
- [18] Kato T, Quasi-linear equations of evolution, with applications to partial differential equations, in: Spectral Theory and Differential Equations, *Lecture Notes in Math.*, Springer-Verlag, Berlin, **448** (1975), 25-70.
- [19] Lenells J, The scattering approach for the Camassa-Holm equation, *J. Nonlinear Math. Phys.* **9** (2002), 389–393.
- [20] Misiolek G, A shallow water equation as a geodesic flow on the Bott-Virasoro group, *J. Geom. Phys.* **24** (1998), 203–208.
- [21] Strichartz R, A Guide to Distribution Theory and Fourier Transforms, CRC Press, Boca Raton, 1994.
- [22] Zhou Y, Blow-up phenomenon for the integrable Degasperis-Procesi equation, *Phys. Lett. A*, in press.
- [23] Yin Z, Global weak solutions for a new periodic integrable equation with peakon solutions, *J. Funct. Anal.* **212** (2004), 182–194.