Study of the Improved Projective Approach and the Variable Separation Solutions in a (2+1)-Dimensional Soliton System

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Abstract. The projective approach is a kind of classic, efficient and well-developed method to solve nonlinear evolution equations, the remarkable characteristic of which is that we can have many different ansatzs and therefore, a large number of solutions. In this paper, with the help of the improved projective approach and a linear variable separation approach, some new variable separation solutions of the (2+1)-dimensional Generalized Calogero-Bogoyavlenskii-Schiff system (GCBS) is derived.

Introduction

Modern soliton theory is widely applied in many natural sciences such as chemistry, biology, mathematics, communication, and in particular in almost all branches of physics like fluid dynamics, plasma physics, field theory, optics, and condensed matter physics [1]. The exact solutions of nonlinear partial differential equations (NPDE) are interesting and popular topic in nonlinear physicists and mathematicians, and various methods for obtaining exact solutions of nonlinear system have been proposed, for example, the bilinear method, the standard Painlevé truncated expansion, the method of “coalescence of eigenvalue” or “wavenumbers”, the homogenous balance method, and the mapping method [2-5] etc. In the past, Mei and Zhang have obtained exact traveling wave solutions for a nonlinear evolution equation with the Riccari equation (\( \xi' = a_0 + a_1 \xi + a_2 \xi^2 \)) projective method [6]. In this paper, by using the Riccari equation projective method, we construct non-traveling wave solutions with

\[ q = lx + my + nt + R(x, y, t) \]

in the (2+1)-dimensional Generalized Calogero-Bogoyavlenskii-Schiff (GCBS) system [7].

\[ \alpha V_{x\tau} + 2 \beta V_x V_{xy} + \beta V_y V_{xx} + V_{xxy} = 0. \]  \hspace{1cm} (1)

where, \( \alpha \) and \( \beta \) are two constants.

As is well known, to search for the solitary wave solutions for a nonlinear physical model, we can apply different approaches. One of the most efficient methods of finding soliton excitations of a physical model is the so-called mapping approach with variable coefficients. The basic ideal of the algorithm is as follows. For a given nonlinear partial differential equation (NPDE) with the independent variables \( x = (x_0, t, x_1, x_2, ..., x_m) \), and the dependent variable \( u \), in the form

\[ P(u, u_t, u_x, u_{x,y}, ...) = 0 \]

(2)

where \( P \) is in general a polynomial function of its arguments, and the subscripts denote the partial derivatives, the solution can be assumed to be in the form,

\[ u = A(x) + \sum_{j=1}^{\infty} (B_j(x) \xi^j[q(x)]) \]

(3)

with

\[ \xi' = a_0 + a_1 \xi + a_2 \xi^2 \]

(4)
and
\[ q = lx + my + nt + R(x, y, t) \]  \hspace{1cm} (5)
where \( a_0, a_1, a_2, l, m, n \) are constants, \( R = R(x, y, t) \) is arbitrary function of \( (x,y,t) \). Substitute (3), (4) and (5) into the given NPDE and collect coefficients of polynomials of \( \xi \), then eliminate each coefficient to derive a set of equations of \( A, B \) and \( q \), and solve these equations to obtain \( A, B \) and \( q \). Finally, as (4) is known to possess the solutions [16], one obtains the complex solutions to the given NPDE.

**New Exact Solutions of the GCBS System**

Now we apply the Riccati equation projective approach to (1). By the balancing procedure, ansatz (3) becomes
\[ V = A(x, y, t) + B(x, y, t)\phi(q(x, y, t)) \]  \hspace{1cm} (6)
where \( A, B, \) and \( q \) are functions of \( (x,y,t) \) to be determined. Substituting (6) and (4) into (1) and collecting coefficients of polynomials of \( \phi \), then setting each coefficient to zero, we have
\[ A = \frac{1}{2\beta} \int \frac{-aq^2 q_y + 4a_0a_2q_y^4 - 4a_1q_y q_{xx} - a_1^2 q_y^4 - 2q_y q_{xxx} + q_{xxx} q_y}{q_y^2} dx, \]
\[ B = -\frac{4a_2q_x}{\beta}, \]  \hspace{1cm} (7)
\[ q = lx + my + nt + R(x, y, t), R = \xi(x) + \varphi(y - cd) \]  \hspace{1cm} (8)

Based on the solutions of (4), we can derive the following complex wave solutions of (1):

**Case 1** when \( a_0=1, a_1=0, a_2=-1 \),
\[ V_1 = \frac{1}{2\beta} \int \frac{-aq^2 q_y + 4a_0a_2q_y^4 - 4a_1q_y q_{xx} - a_1^2 q_y^4 - 2q_y q_{xxx} + q_{xxx} q_y}{q_y^2} dx \]
\[ \frac{4a_2q_x}{\beta} \tanh(q), \]  \hspace{1cm} (9)
\[ V_2 = \frac{1}{2\beta} \int \frac{-aq^2 q_y + 4a_0a_2q_y^4 - 4a_1q_y q_{xx} - a_1^2 q_y^4 - 2q_y q_{xxx} + q_{xxx} q_y}{q_y^2} dx \]
\[ \frac{4a_2q_x}{\beta} \coth(q), \]  \hspace{1cm} (10)

**Case 2** when \( a_0=1, a_1=0, a_2=1 \),
\[ V_3 = \frac{1}{2\beta} \int \frac{-aq^2 q_y + 4a_0a_2q_y^4 - 4a_1q_y q_{xx} - a_1^2 q_y^4 - 2q_y q_{xxx} + q_{xxx} q_y}{q_y^2} dx \]
\[ \frac{4a_2q_x}{\beta} \tan(q), \]  \hspace{1cm} (11)

**Case 3** when \( a_0=-1, a_1=0, a_2=-1 \),
\[ V_4 = \frac{1}{2\beta} \int \frac{-aq^2 q_y + 4a_0a_2q_y^4 - 4a_1q_y q_{xx} - a_1^2 q_y^4 - 2q_y q_{xxx} + q_{xxx} q_y}{q_y^2} dx \]
\[ \frac{4a_2q_x}{\beta} \cot(q), \]  \hspace{1cm} .
\[ V_s = \frac{1}{2\beta} \int -aq^2g + 4a_0a_0q^4g - 4a_1q, q_x, q_x - a_2q^4, q_y, -2q, q_x, q_x + q_x^3, q_y \, dx \]
\[ -\frac{4a_2q_x}{\beta} (\tan(q) + \sec(q)), \]  
(13)

\[ V_5 = \frac{1}{2\beta} \int -aq^2g + 4a_0a_0q^4g - 4a_1q, q_x, q_x - a_2q^4, q_y, -2q, q_x, q_x + q_x^3, q_y \, dx \]
\[ -\frac{4a_2q_x}{\beta} (\tan(q) - \sec(q)), \]  
(14)

\[ V_7 = \frac{1}{2\beta} \int -aq^2g + 4a_0a_0q^4g - 4a_1q, q_x, q_x - a_2q^4, q_y, -2q, q_x, q_x + q_x^3, q_y \, dx \]
\[ -\frac{4a_2q_x}{\beta} (\csc(q) - \cot(q)), \]  
(15)

**Case 5 when** \( a_0 = -\frac{1}{2}, a_1 = 0, a_2 = -\frac{1}{2} \),

\[ V_9 = \frac{1}{2\beta} \int -aq^2g + 4a_0a_0q^4g - 4a_1q, q_x, q_x - a_2q^4, q_y, -2q, q_x, q_x + q_x^3, q_y \, dx \]
\[ -\frac{4a_2q_x}{\beta} (-\tan(q) + \sec(q)), \]  
(16)

\[ V_{10} = \frac{1}{2\beta} \int -aq^2g + 4a_0a_0q^4g - 4a_1q, q_x, q_x - a_2q^4, q_y, -2q, q_x, q_x + q_x^3, q_y \, dx \]
\[ -\frac{4a_2q_x}{\beta} (\cot(n) + \csc(q)), \]  
(17)

\[ V_{11} = \frac{1}{2\beta} \int -aq^2g + 4a_0a_0q^4g - 4a_1q, q_x, q_x - a_2q^4, q_y, -2q, q_x, q_x + q_x^3, q_y \, dx \]
\[ -\frac{4a_2q_x}{\beta} (\cot(n) - \csc(q)), \]  
(18)

\[ V_{12} = \frac{1}{2\beta} \int -aq^2g + 4a_0a_0q^4g - 4a_1q, q_x, q_x - a_2q^4, q_y, -2q, q_x, q_x + q_x^3, q_y \, dx \]
\[ -\frac{4a_2q_x}{\beta} (\frac{\cot(q)}{1 + \csc(q)}), \]  
(19)

\[ V_{13} = \frac{1}{2\beta} \int -aq^2g + 4a_0a_0q^4g - 4a_1q, q_x, q_x - a_2q^4, q_y, -2q, q_x, q_x + q_x^3, q_y \, dx \]
\[ -\frac{4a_2q_x}{\beta} (\frac{\cot(q)}{1 - \csc(q)}), \]  
(20)

**Case 6 when** \( a_0 = \frac{1}{2}, a_1 = 0, a_2 = -\frac{1}{2} \)
\[ V_{13} = \frac{1}{2\beta} \int -a_0^2 q_i + 4a_0 a_i q_i^4 q_y - 4a_i q_i q_y q_{xy} - a_i^2 q_i^4 q_y - 2q_i q_y q_{xy} + q_y q_{xy} \, dx - \frac{4a_i q_y}{\beta} \]

(coth(q) + \csc h(q)),

(21)

\[ V_{14} = \frac{1}{2\beta} \int -a_0^2 q_i + 4a_0 a_2 q_i^4 q_y - 4a_2 q_i q_y q_{xy} - a_2^2 q_i^4 q_y - 2q_i q_y q_{xy} + q_y q_{xy} \, dx - \frac{4a_2 q_y}{\beta} \]

- (coth(q) - \csc h(q)),

(22)

\[ V_{15} = \frac{1}{2\beta} \int -a_0^2 q_i + 4a_0 a_2 q_i^4 q_y - 4a_2 q_i q_y q_{xy} - a_2^2 q_i^4 q_y - 2q_i q_y q_{xy} + q_y q_{xy} \, dx - \frac{4a_2 q_y}{\beta} \]

- (tan(q) + I sec h(q)),

(23)

\[ V_{16} = \frac{1}{2\beta} \int -a_0^2 q_i + 4a_0 a_2 q_i^4 q_y - 4a_2 q_i q_y q_{xy} - a_2^2 q_i^4 q_y - 2q_i q_y q_{xy} + q_y q_{xy} \, dx - \frac{4a_2 q_y}{\beta} \]

- (tan(q) - I sec h(q)),

(24)

\[ V_{17} = \frac{1}{2\beta} \int -a_0^2 q_i + 4a_0 a_2 q_i^4 q_y - 4a_2 q_i q_y q_{xy} - a_2^2 q_i^4 q_y - 2q_i q_y q_{xy} + q_y q_{xy} \, dx - \frac{4a_2 q_y}{\beta} \]

- tan(q)

(25)

\[ V_{18} = \frac{1}{2\beta} \int -a_0^2 q_i + 4a_0 a_2 q_i^4 q_y - 4a_2 q_i q_y q_{xy} - a_2^2 q_i^4 q_y - 2q_i q_y q_{xy} + q_y q_{xy} \, dx - \frac{4a_2 q_y}{\beta} \]

- tan(q)

(26)

\[ V_{19} = \frac{1}{2\beta} \int -a_0^2 q_i + 4a_0 a_2 q_i^4 q_y - 4a_2 q_i q_y q_{xy} - a_2^2 q_i^4 q_y - 2q_i q_y q_{xy} + q_y q_{xy} \, dx - \frac{4a_2 q_y}{\beta} \]

- tan(q)

(27)

\[ V_{20} = \frac{1}{2\beta} \int -a_0^2 q_i + 4a_0 a_2 q_i^4 q_y - 4a_2 q_i q_y q_{xy} - a_2^2 q_i^4 q_y - 2q_i q_y q_{xy} + q_y q_{xy} \, dx - \frac{4a_2 q_y}{\beta} \]

- tan(q)

(28)

\[ \text{Case 7. when } a_0 = 1, a_1 = 0, a_2 = -4, \]
\[ V_{21} = \frac{1}{2\beta} \int \left( -a q_x^2 q_y^2 + 4a_0 a_2 q_x^4 q_y - 4a_1 q_x q_y q_x q_xx - a_1^2 q_x^4 q_y - 2q_x q_y q_x q_xx + q_x q_y \right) dx \\
- \frac{4a_2 q_x}{\beta} \frac{\tanh(q)}{1 + \tanh^2(q)} \right], \\
(29) \\
Case 8. when \(a_0 = 1, a_1 = 0, a_2 = 4, \)
\[ V_{22} = \frac{1}{2\beta} \int \left( -a q_x^2 q_y^2 + 4a_0 a_2 q_x^4 q_y - 4a_1 q_x q_y q_x q_xx - a_1^2 q_x^4 q_y - 2q_x q_y q_x q_xx + q_x q_y \right) dx \\
- \frac{4a_2 q_x}{\beta} \frac{\tan(q)}{1 - \tan^2(q)} \right], \\
(30) \\
Case 9. when \(a_0 = -1, a_1 = 0, a_2 = -4, \)
\[ V_{23} = \frac{1}{2\beta} \int \left( -a q_x^2 q_y^2 + 4a_0 a_2 q_x^4 q_y - 4a_1 q_x q_y q_x q_xx - a_1^2 q_x^4 q_y - 2q_x q_y q_x q_xx + q_x q_y \right) dx \\
- \frac{4a_2 q_x}{\beta} \frac{\cot(q)}{1 - \cot^2(q)} \right], \\
(31) \\
Case 10. when \(a_0 = 1, a_1 = 2, a_2 = 2, \)
\[ V_{24} = \frac{1}{2\beta} \int \left( -a q_x^2 q_y^2 + 4a_0 a_2 q_x^4 q_y - 4a_1 q_x q_y q_x q_xx - a_1^2 q_x^4 q_y - 2q_x q_y q_x q_xx + q_x q_y \right) dx \\
- \frac{4a_2 q_x}{\beta} \frac{\tan(q)}{1 - \tan(q)} \right], \\
(32) \\
Case 11. when \(a_0 = 1, a_1 = -2, a_2 = 2, \)
\[ V_{25} = \frac{1}{2\beta} \int \left( -a q_x^2 q_y^2 + 4a_0 a_2 q_x^4 q_y - 4a_1 q_x q_y q_x q_xx - a_1^2 q_x^4 q_y - 2q_x q_y q_x q_xx + q_x q_y \right) dx \\
- \frac{4a_2 q_x}{\beta} \frac{\tan(q)}{1 + \tan(q)} \right], \\
(33) \\
Case 12. when \(a_0 = -1, a_1 = 2, a_2 = -2, \)
\[ V_{26} = \frac{1}{2\beta} \int \left( -a q_x^2 q_y^2 + 4a_0 a_2 q_x^4 q_y - 4a_1 q_x q_y q_x q_xx - a_1^2 q_x^4 q_y - 2q_x q_y q_x q_xx + q_x q_y \right) dx \\
- \frac{4a_2 q_x}{\beta} \frac{\cot(q)}{1 + \cot(q)} \right], \\
(34) \\
Case 13. when \(a_0 = -1, a_1 = -2, a_2 = 2, \)
\[ V_{27} = \frac{1}{2\beta} \int \left( -a q_x^2 q_y^2 + 4a_0 a_2 q_x^4 q_y - 4a_1 q_x q_y q_x q_xx - a_1^2 q_x^4 q_y - 2q_x q_y q_x q_xx + q_x q_y \right) dx \\
- \frac{4a_2 q_x}{\beta} \frac{\cot(q)}{1 - \cot(q)} \right], \\
(35) \\
Case 14. when \(a_0 = 0, a_1 = 0, a_2 = a, \)
\[ V_{2s} = \frac{1}{2\beta} \int \left[ -\alpha q_x^2 q_y + 4a_1 a_2 q_x^2 q_y - 4a_1 q_x q_y q_{xx} - a_1^2 q_x^4 q_y - 2a_1 q_x q_y q_{xxy} + q_x q_y \right] dx \]

\[ + \frac{4a_2 q_x}{\beta} \frac{1}{aq + k}, \]

(36)

with \( q = lx + my + nt + R, \ R = \chi(x) + \varphi(y - ct) \). Here, \( \alpha, \beta, c, a \) and \( k \) are constants.

**Conclusions**

The projective equation method is a kind of classic, efficient and well-developed method to solve nonlinear evolution equations. In the past, Mei and Zhang have obtained the exact traveling wave solutions for the nonlinear evolution equations such as Gross-Pitaevskii equation with the Riccari equation projective approach. In this paper, we extend this approach and construct the variable separation solutions of the (2+1)-dimensional Generalized Calogero-Bogoyavlenskii-Schiff system, which are different from the ones of the previous work. Since the wide applications of the soliton theory, to learn more about the localized excitations and their applications in reality is worthy of study further.

**References**


