

A note on Bernoulli polynomials and solitons

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Abstract

The dependence on time of the moments of the one-soliton KdV solutions is given by Bernoulli polynomials. Namely, we prove the formula

$$\int_{\mathbb{R}} x^n \operatorname{sech}^2(x-t) dx = 2\pi^n (-i)^n B_n\left(\frac{1}{2} + \frac{t}{\pi}i\right),$$

expressing the moments of the one-soliton function $\operatorname{sech}^2(x-t)$ in terms of the Bernoulli polynomials $B_n(x)$. We also provide an alternative short proof to the Grosset-Veselov formula connecting the one-soliton to the Bernoulli numbers

$$\int_{\mathbb{R}} (D^{m-1} \operatorname{sech}^2 x)^2 dx = (-1)^{m-1} 2^{2m+1} B_{2m},$$

($D = d/dx$) published recently in this journal.

1 Introduction

The Bernoulli polynomials $B_n(t)$ are defined by the generating function

$$\frac{z e^{zt}}{e^z - 1} = \sum_{n=0}^{\infty} B_n(t) \frac{z^n}{n!}, \quad (1.1)$$

($|z| < 2\pi$) and the Bernoulli numbers are their values at zero, $B_n = B_n(0)$. In two recent papers Fairlie and Veselov [2] and Grosset and Veselov [4] revealed an interesting connection between Bernoulli polynomials and the theory of the Korteweg-deVries equation

$$u_t - 6u u_x + u_{xxx} = 0, \quad (1.2)$$

and in particular, with its remarkable single soliton solution [5]

$$u(x, t) = -2 \operatorname{sech}^2(x - 4t). \quad (1.3)$$

In their letter [4] Grosset and Veselov established the formula

$$B_{2m} = \frac{(-1)^{m-1}}{2^{2m+1}} \int_{\mathbb{R}} (D^{m-1} \operatorname{sech}^2 x)^2 dx, \quad (1.4)$$

which connects the Bernoulli numbers and the derivatives of the single soliton. In section 3 we give a short proof of this formula based on Fourier theory. First of all, we point out another interesting connection between the single soliton and the Bernoulli polynomials.

2 Main result

Writing $\operatorname{sech}(x) = 1/\cosh(x)$, we have the following proposition.

Proposition 1. For all t and $n = 0, 1, 2, \dots$,

$$\int_{\mathbb{R}} \frac{x^n}{\cosh^2(x-t)} dx = 2(-i\pi)^n B_n\left(\frac{1}{2} + \frac{t}{\pi}i\right). \quad (2.1)$$

Remark 1. The Bernoulli polynomials have the addition property [7, p. 4]

$$B_n(x+y) = \sum_{k=0}^n \binom{n}{k} B_k(x) y^{n-k} \quad (2.2)$$

and also the property

$$B_k\left(\frac{1}{2}\right) = (2^{1-k} - 1) B_k. \quad (2.3)$$

Therefore, we can put (2.1) in the form

$$\int_{\mathbb{R}} \frac{x^n}{\cosh^2(x-t)} dx = 2 \sum_{k=0}^n \binom{n}{k} (-\pi i)^k (2^{1-k} - 1) t^{n-k} B_k, \quad (2.4)$$

where on the right hand side all terms with odd indices are zeros, as $B_k = 0$ when $k > 1$ is odd, and for $k = 1$ the factor $2^{1-k} - 1$ is zero.

Remark 2. If

$$f(x) = \sum_{n=0}^m a_n x^n \quad (2.5)$$

is a polynomial, we can multiply equation (2.1) by a_n and sum for $n = 0, 1, \dots, m$ to obtain the superposition formula

$$\int_{\mathbb{R}} \frac{f(x)}{\cosh^2(x-t)} dx = 2 \sum_{n=0}^m a_n (-i\pi)^n B_n\left(\frac{1}{2} + \frac{t}{\pi}i\right). \quad (2.6)$$

Proof of the Proposition. We use the Fourier transform formula [3, 3.982.1, p. 505] or [6, 1.7.2, p. 33]

$$\int_{\mathbb{R}} \frac{e^{ixy}}{\cosh^2 x} dx = \frac{\pi y}{\sinh \frac{\pi y}{2}}, \quad (2.7)$$

and change the variable $x \rightarrow x - t$ to obtain

$$\int_{\mathbb{R}} \frac{e^{ixy}}{\cosh^2(x-t)} dx = \frac{\pi y e^{ity}}{\sinh \frac{\pi y}{2}} = \frac{2\pi y e^{\pi y(\frac{1}{2} + \frac{it}{\pi})}}{e^{\pi y} - 1}. \quad (2.8)$$

This is explicitly the Fourier transform of the single soliton solution. Next, we expand the right hand side of (2.8) in a power series on the powers of πy for $|y| < 2$. In view of (1.1) this provides the representation

$$\int_{\mathbb{R}} \frac{e^{ixy}}{\cosh^2(x-t)} dx = 2 \sum_{n=0}^{\infty} B_n \left(\frac{1}{2} + \frac{t}{\pi} i\right) \frac{\pi^n y^n}{n!}. \quad (2.9)$$

Expanding now in power series the exponential function inside the integral, changing the order of summation and integration, and comparing the coefficients for y^n we arrive at (2.1). The proof is complete. \blacksquare

3 The Grosset-Veselov Formula

We show here that Grosset-Veselov's formula (1.4) is equivalent to the representation

$$B_{2m} = \frac{(-1)^{m-1}}{\pi^{2m}} \int_0^{+\infty} \frac{x^{2m}}{\sinh^2 x} dx, \quad (3.1)$$

via Fourier transform theory. Formula (3.1) is known and can be found in [1, 1.13 (27)] or [3, 3.527.2, p. 352]. To understand the nature of (3.1) better, we give a short derivation in the Appendix.

Proof. We use the notation

$$F(t) = \int_{\mathbb{R}} f(x) e^{-2\pi i x t} dx \quad (3.2)$$

for the Fourier transform and write (2.7) as

$$\int_{\mathbb{R}} \frac{e^{-2\pi i x t}}{\cosh^2 x} dx = \frac{2\pi^2 t}{\sinh(\pi^2 t)}. \quad (3.3)$$

According to the derivative property we find

$$\int_{\mathbb{R}} \left(D^{m-1} \frac{1}{\cosh^2 x} \right) e^{-2\pi i x t} dx = (2\pi i t)^{m-1} \frac{2\pi^2 t}{\sinh(\pi^2 t)} = \frac{-\pi i (2\pi i t)^m}{\sinh(\pi^2 t)}. \quad (3.4)$$

Applying now Plancherel's theorem, i.e.

$$\int_{\mathbb{R}} |F(t)|^2 dt = \int_{\mathbb{R}} |f(x)|^2 dx, \quad (3.5)$$

we obtain

$$\int_{\mathbb{R}} \left(D^{m-1} \frac{1}{\cosh^2 x} \right)^2 dx = \int_{\mathbb{R}} \left| \frac{\pi i (2\pi i t)^m}{\sinh(\pi^2 t)} \right|^2 dt = 2^{2m+1} \pi^{2m+2} \int_0^{+\infty} \frac{t^{2m}}{\sinh^2(\pi^2 t)} dt. \quad (3.6)$$

Therefore, according to (3.1),

$$\int_{\mathbb{R}} \left(D^{m-1} \frac{1}{\cosh^2 x} \right)^2 dx = (-1)^{m-1} 2^{2m+1} B_{2m}, \quad (3.7)$$

which is Grosset-Veselov's formula. \blacksquare

We want to point out that this proof was independently suggested by Professor A. Staruszkiewicz (see Note added in Proofs at the end of [4]). For the convenience of the reader it is appropriate to have it recorded here together with the derivation of (3.1) below.

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Appendix

A

We deduce (3.1) from the well-known Fourier sine transform formula [3, 3.911.2, p. 481] or [6, 2.3.12, p. 125]

$$\int_0^{+\infty} \frac{\sin(ty)}{e^t - 1} dt = \frac{\pi}{2} \coth(\pi y) - \frac{1}{2y}. \quad (\text{A.1})$$

After integration by parts the left hand side takes the form

$$\frac{1}{y} \int_0^{+\infty} \frac{(1 - \cos(ty)) e^t}{(e^t - 1)^2} dt = \frac{1}{2y} \int_0^{+\infty} \frac{1 - \cos(2xy)}{\sinh^2 x} dx \quad (\text{setting } t = 2x). \quad (\text{A.2})$$

Equation (A.1) can now be written as

$$\int_0^{+\infty} \frac{1 - \cos(2xy)}{\sinh^2 x} dx = \pi y - 1 + \frac{2\pi y}{e^{2\pi y} - 1}. \quad (\text{A.3})$$

Expanding both sides on powers of y by using the two series

$$1 - \cos(2xy) = \sum_{n=1}^{+\infty} \frac{(-1)^{n-1} 2^{2n} x^{2n} y^{2n}}{(2n)!}; \quad (\text{A.4})$$

$$\pi y - 1 + \frac{2\pi y}{e^{2\pi y} - 1} = \sum_{n=1}^{+\infty} B_{2n} \frac{(2\pi)^{2n} y^{2n}}{(2n)!}, \quad (\text{A.5})$$

and comparing coefficients we arrive at (3.1).

References

- [1] ERDÉLYI A (editor), Higher Transcendental Functions, vol. 1, McGraw-Hill, New York, 1955.
- [2] FAIRLIE D B and VESELOV A P, Faulhaber and Bernoulli polynomials and solitons, *Physica D* **152-153** (2001), 47–50.
- [3] GRADSHTEYN I S and RYZHIK I M, Tables of Integrals, Series and Products, Academic Press, London, 1980.

- [4] GROSSET M-P and VESELOV A P, Bernoulli polynomials and solitons, *J. Nonlinear Math. Phys.* **12** (4) (2005), 469–474.
- [5] LAMB G L JR, *Elements of Soliton Theory*, John Wiley, New York, 1980.
- [6] OBERHETTINGER O, *Tables of Fourier Transforms*, Springer Verlag, Berlin, 1990.
- [7] TEMME N M, *Special Functions*, John Wiley, New York, 1996.